TCC-III Engine Benchmark for Large-Eddy Simulation of IC Engine Flows

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Abstract — A collaborative effort is described to benchmark the TCC-III engine, and to illustrate the application of this data for the evaluation of sub-grid scale models and valve simulation details on the fidelity of Large-Eddy Simulations (LES). The TCC-III is a spark ignition 4-stroke 2-valve engine with a flat head and piston and is equipped with a full quartz liner for maximum optical access that allows high-speed flow measurements with Particle Image Velocimetry (PIV); the TCC-III has new valve seats and a modified intake-system compared to previous configurations. This work is an extension of a previous study at an engine speed of 800 RPM and an intake manifold pressure (MAP) of 95 kPa, where a one-equation eddy viscosity LES model yielded accurate qualitative and quantitative predictions of ensemble averaged mean and RMS velocities during the intake and compression stroke. Here, experimental data were acquired with parametric variation of engine speed and intake manifold absolute pressure to assess the capability of LES models over a range of operating conditions of practical relevance. This paper focuses on the repeatability and accuracy of the measured PIV data, acquired at 1 300 RPM, at two different MAP (95 kPa and 40 kPa), and imaged at multiple data planes and crank angles. Two examples are provided, illustrating the application of this data to LES model development. In one example, the experimental data are used to distinguish between the efficacies of a one-equation eddy viscosity model versus a dynamic structure one-equation model for the sub-grid stresses. The second example addresses the effects of numerical intake-valve opening strategy and local mesh refinement in the valve curtain.

Résumé — Benchmark de moteur de référence TCC-III pour la simulation aux grandes échelles (Large-Eddy Simulations, LES) de l’écoulement dans les moteurs à combustion interne — Un projet collaboratif est décrit, visant à caractériser le moteur TCC-III et à illustrer l’application de ces données pour l’évaluation de modèles de sous-maille et des détails de simulation des soupapes sur la fidélité de simulations aux grandes échelles. Le TCC-III est un moteur à allumage commandé 4 temps à deux soupapes et à culasse et piston plats, et il est équipé d’un cylindre en quartz donnant un accès optique maximal qui permet des mesures par vélocimétrie par image de particules (Particle Image Velocimetry, PIV) rapide ; le TCC-III présente de nouveaux sièges de soupapes et un système d’admission modifié par rapport aux configurations précédentes. Le présent travail représente une extension d’une étude précédente à un régime moteur de 800 tr/min et à une pression de collecteur d’admission de 95 kPa, pour laquelle un modèle LES de viscosité turbulente à une équation avait donné des prédictions qualitatives et quantitatives précises de la moyenne d’ensemble du champ de vitesse et de sa fluctuation pendant l’admission et la phase de compression. Dans le présent travail,
les mesures expérimentales ont été acquises avec une variation paramétrique du régime moteur et de la pression absolue du collecteur d’admission, afin d’évaluer la capacité des modèles LES sur une gamme de conditions opératoires d’intérêt pratique. Le présent article met l’accent sur la répétabilité et la précision des mesures PIV, acquises à 1 300 tr/min, à deux pressions collecteur admission différentes (95 kPa et 40 kPa), et réalisées dans plusieurs plans de visualisation et à différents angles vilebrequin. Deux exemples sont fournis, illustrant l’application de ces données au développement du modèle LES. Dans un exemple, les données expérimentales sont utilisées pour distinguer les efficacités d’un modèle de viscosité turbulente à une équation par rapport à un modèle de structure dynamique à une équation pour les contraintes sous-maille. Le second exemple traite des effets de la stratégie numérique d’ouverture de la soupape d’admission et du raffinement de maillage local dans le rideau de soupapes.

NOMENCLATURE

INTRODUCTION

Reynolds-Average Navier Stokes (RANS) simulation of reciprocating Internal Combustion Engine (ICE) combustion has matured to be a valuable tool for engineering design. This has been especially true for evaluating the average-cycle behavior of new engine concepts [1]. With increased computational capabilities, more details, such as cyclic variability (Cycle-to-Cycle Variation, CCV) can be addressed using multi-cycle Large-Eddy Simulation (LES), that has been evolving since the 1990s [2]. Since then, LES has been used to simulate flow and combustion in modern engine concepts such as four-valve pent-roof spark-ignition, Diesel, and HCCI engines [3-9]. These simulations of comprehensive and complex engine architecture are useful not only to guide engine development but also to identify the needed improvements in the LES approaches. In particular, for computational efficiency and cost considerations these simulation approaches must employ practical meshes and make simplifying assumptions, which include sub-models to describe processes below the resolved spatial or temporal scales. The accuracy of the sub-grid-models (Sub-Grid-Scale, SGS), the numerical schemes in use, and more, have an impact on how well the in-cylinder processes can be simulated. These computations thus rely on accurate experimental data to identify the underlying physics and benchmark the simulation results. Only with accurate and repeatable experimental data is it possible to achieve a meaningful assessment and validation of the simulations. With this purpose in mind, the Transparent Combustion Chamber (TCC-0) engine was originally designed (circa 1990), built, and used for more fundamental investigations of the in-cylinder flow and combustion CCV using the then nascent Particle Image Velocimetry (PIV). The engine design was intended to provide benchmark data at the next logical level beyond the Imperial College data [10], adding essential features of ICE operation such as complete charge transfer processes, continuous operation, and combustion. Complementing the more recent LES calculations noted above, there are
The TCC engine was resurrected in 2010 in the TCC-II version, to renew the fundamental investigation of CCV using experiments in conjunction with the evaluation of three different LES approaches [13]. The rationale to revert to this engine configuration was based on the original design philosophy. In particular, the geometry was purposefully chosen to be simple for ease of computational gridding, and optimized for optical access to the entire combustion chamber. In addition, the two-valve, pancake chamber design with a large piston-diameter to valve-diameter ratio produces extremely large-scale flow CCV when compared to the more naturally directed flow of a 4-valve pent roof design. The intent was to create an extremely large-scale flow CCV that would be obviously detectable and an extreme test for simulations. It was used in its original configuration, the TCC-0, for experiments prior to 2000 [14-17]. The TCC-II configuration had the same engine geometry (refer to Tab. A1 of Appendix A) but a different intake system, exhaust system, and valve seat (two angles) compared to the TCC-0. The experimental results from the TCC-II configuration were useful for establishing the LES-to-experimental data-comparison protocols, to identify desired improvements in the hardware and operating procedures, and to assess the impact of engine and system imperfections on the in-cylinder flow [18-20]. Based on the TCC-II studies, new results are reported here from a third configuration, TCC-III, which has refurbished valve hardware and the intake/exhaust systems upgraded for fired testing. Also, the TCC-III data are significantly more repeatable and with expanded operating conditions as listed in Tables A2 and A3, respectively. Chronicles of the TCC engine versions and cross references to the corresponding publications can be found under http://deepblue.lib.umich.edu/handle/2027.42/108382.

The purpose of this paper is twofold. First, to document the motored experimental data and, second, to show examples of its use to guide the setup of LES and selection of sub-models to best capture experimentally observed flow conditions. To this end, the accuracy and test-to-test repeatability of the intra-cycle, high-speed motored pressure and PIV velocity data is documented. Then it is demonstrated how the data is used to make informed selections of LES SGS models and local mesh refinement efforts in the valve curtain. Given their frequent use, an eddy viscosity model and a dynamic structure model were chosen here for illustration. It is important to assess simulation performance for a range of operating conditions to avoid optimization or tuning of simulation settings to one condition but not gaining insight into the fidelity of reproducing different conditions. Two operating conditions will be highlighted here to illustrate how changes in intake manifold pressure affect the filling dynamics of the cylinder and how the simulations pick up this process. Conceptually, it is expected that a dynamic structure-based subgrid treatment would perform better than a simple eddy viscosity model in handling the flow around the valves and the resulting strong velocity gradients [8] in the intake ‘jet’ flow. Encouraging results that point to this led to investigations of increasing the mesh resolution around the valves to better capture the dynamics of the valve opening and closing events; this resulted in higher accuracy for the digital nature of valving events in CFD.

The experimental flow and pressure data acquired in this study is available upon request, while the geometry of the TCC-III engine and the GT Power model can be downloaded at http://deepblue.lib.umich.edu/handle/2027.42/108382 with the purpose of helping the engine CFD community to develop more accurate flow models.

1 EXPERIMENTAL SETUP AND PROCEDURES

1.1 Engine Systems

The UM (University of Michigan) TCC-III optical engine used here is a spark-ignition 2-valve, 4-stroke, pancake shape combustion chamber engine with a geometrical compression ratio of 10:1 and bore × stroke of 92 × 86 mm (Fig. 1). It is equipped with a full quartz cylinder and a 70 mm-diameter flat quartz piston window. Geometry and valve timings are provided in Table A1.

The TCC-III configuration has hardware upgrades from the TCC-II configuration. The TCC-III engine was refurbished with new hydraulic valve lifters, new valves (identical design), and new four-angle valve seats. Optical methods were established to determine the valve opening and closing times. The improvements were made in response to TCC-II simulation and experimental studies that revealed the need for more repeatable and better-documented valve events. The intake and exhaust systems are now prepared for firing tests as shown schematically in Figure 2a. Both a gaseous-fuel and a nitrogen-dilution critical-orifice metering systems have been added. Flame arrestors have been added at the plenum inlets/outlets and water-cooled exhaust-pressure transducers were installed. The fuel and dilution systems were connected during the motored tests presented here, in order to sustain the same intake/exhaust system wave dynamics with fired operation (not presented here). Thus, new engine geometry files are required both for the LES and 1-D (GT Power) simulations of the TCC-III engine. Detailed geometry, .stl, .igs, and the full GT Power model along with large experimental data sets (pressure and velocity) are

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1 The TCC-I configuration did not exist in hardware but was a simulation-only geometry identical to the TCC-II hardware, except for utilizing a 45-degree, single-angle valve seat.
The intake and exhaust systems were instrumented with high-frequency piezo-resistive absolute-pressure transducers (Kistler 4007B in the intake and Kistler 4049A in the exhaust (water cooled)) to record the intra-cycle pressure every 0.5 CAD and to determine CCV at the desired boundary conditions (Fig. 1). The intake plenum and runner, exhaust runner, and cooling water were electrically heated to the 45°C intake-air operating point to minimize temperature gradients in the air flow. The engine was warmed up and transducer zero-drift adjusted to barometric standards prior to each test. The uncertainty of the transducers is estimated to be 0.1-0.2% of full scale output = ±0.5-1 kPa for the 500 kPa transducers in use. A back-pressure valve just downstream of the flame arrester was used to maintain the exhaust plenum operation set point at 101.5 kPa average pressure.

Operating procedures were instituted to assure engine operation repeatability during the six months the data was acquired, including charting of the run-time control parameters. The engine was motored for 15-20 minutes to achieve conditions that were nearly in steady-state, and repeatable. The temperature of the quartz cylinder outside-surface at mid-stroke was brought to 40±4°C, depending on operating condition. At run time, the engine speed, delivered air mass flow, and the average intake and exhaust pressures were controlled. The average port pressures, peak cylinder pressure and IMEP were monitored to ensure the engine trapped mass was repeating. Table A2 gives an overview of test-to-test variation of several critical engine parameters and their stability throughout each run.

The test-to-test repeatability of the engine operation is illustrated in Figure 3a, where differences between the individual-test average (over hundreds of cycles) and the average of all tests (discrepancy) are quantified. Here, the discrepancy is plotted since the differences are too small to discriminate on a full-scale plot of the absolute values. The test-to-test deviation in the plenum and ports is much less than 0.5%. In-cylinder pressures are similar through much of the cycle but deviate up to 2% during the valve overlap and early part of the intake stroke (−45 to 90 aTDCexh). At 2% these discrepancies are still small compared to differences shown between the simulations and the data (Fig. 3b), and any discussion of their cause would be pure speculation at this time. The percent discrepancy for the 40 kPa operation is higher due to the lower absolute value, but nominally is of the same order as the 95 kPa operation. Note that data set names are included in many figures throughout this paper to facilitate the connection with the shared data sets (Tab. A3).

1.2 PIV Measurements

The PIV measurements were recorded successively in the four planes shown in Figure 4, every 5 CAD for at least 235 consecutive cycles per test. The Field Of View (FOV) is restricted by the piston window to approximately the center 70 mm for both vertical and horizontal cutting planes. In both the \(z = -5\) mm and \(z = -30\) mm planes the light sheet is brought into the cylinder from the positive \(x\)-direction. Thus, during intake stroke a major part of the FOV is blocked by the intake valve and its shade for the \(z = -5\) mm images. The vertical cutting planes have a FOV from piston to 1% at 95 kPa MAP. The piezo-electric cylinder pressure transducer was pegged using a 10 CAD average at 140 aTDCexh, where the 1-D model predicted that pressures in the intake plenum, the intake port, and in the cylinder should all be at the same value. A back-pressure valve just downstream of the flame arrester was used to maintain the exhaust plenum operation set point at 101.5 kPa average pressure.

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cylinder head with an approximate width of just below 70 mm. A summary is provided in Table A3. Images were recorded using a monochrome high-speed camera (Vision Research, Phantom v1610) with a 1280 × 980 pixel sensor and 12-bit dynamic range. Silicone-oil droplets (1 μm) were added to the intake air and illuminated by a high-repetition-rate dual-cavity frequency doubled Nd:YLF laser (Darwin Duo, Quantronix). The light sheet thickness was 2.0-2.4 mm in the various cutting planes adjusted based on the in-plane resolution set by the 210 mm Nikon Micro-Nikkor ED lens with an aperture of f/# = 5.6 and the distance to the imaging plane.

Image processing began by remapping the raw images to eliminate the non-linear grid distortions caused by the cylinder and thick piston window. A 3rd-order polynomial fit was used for all data recorded in a vertical plane and a linear fit for data captured in the horizontal planes. Then a sliding minimum intensity subtraction was performed to reduce reflections from cylinder head, valve edges and spark plug. All vector fields were calculated using a commercial PIV code (DaVis 8.x, LaVision) employing a decreasing interrogation window size (1 × 128 × 128 pixel, 50% overlap then 2 × 32 × 32 pixel, 50% overlap). The final interrogation window spot size is 2.5 to 2.8 mm with a vector separation of 1.25-1.4 mm. The percentage of first-choice vectors of the experimental data shown here is >98%, meaning that less than 2% of the vectors of the instantaneous flow fields have been lower than 1st choice or have been interpolated.
Experimental pressure test-to-test repeatability compared to simulation-to-measured pressures. a) Three test average pressure traces (in kPaA = kPa absolute) and discrepancy of the individual test towards this average (and GT Power for plenums). b) Discrepancy of simulations at 40 and 95 kPa MAP for port and cylinder pressures.

Figure 3
velocity error (resolution) and maximum velocity as a function of crank angle based on the 0.2 pixel detection limit and 8-pixel maximum, respectively. Figure 5 can then be used to estimate the experimental error, which is important when making quantitative comparisons with simulation data.

Experience tells that measurements in engines can show substantial test-to-test variability and therefore it is essential for LES validation to quantify the test-to-test repeatability of the velocity measurements. For this assessment and to illustrate this here, the ensemble average velocity was sampled along the centerline of the cylinder \((x = 0, y = 0)\) for the \(x-z\) and \(y-z\) planes as shown in Figure 6. The planes required different optical set-ups, and the three tests were taken at early, middle, and late times during the six-month measurement campaign. As with the pressure, the velocity repeatability will be shown to be much better than the discrepancies observed between experiments and simulations.

The quality of convergence of the ensemble-averaged velocity and ensemble-standard deviation of velocities are a function of the number of cycles used to compute the values, the location, and crank angle. To illustrate this, the coefficients of variance \(COV\) at 100 and 300 CAD aTDC are shown in Figure 7a, which can be used as a metric to estimate the deviation of the calculated mean flow to the converged mean. Regions of low \(COV\) are expected to converge faster than high \(COV\) areas. For a more in-depth discussion the convergence at point A is examined in greater detail and those results are shown in Figure 7b-d. In Figure 7b, the normalized ensemble average velocity magnitude of different-sized sub-samples is shown for consecutive and randomized cycles. Using randomized sub-samples is an important statistical check since it can be assumed that there might be some correlation in flows from one cycle to the next. In particular, due to slow drifts in the mean flow with time (shown in Fig. 7d), caused by changes in swirl ratio [19], the variations are interdependent. The solid black lines show the estimated error in the mean that was determined as sample size divided by \(COV\) squared as proposed by [24] and indicate about the 80% confidence interval. Doubling the error range (dashed lines) shows about the entire range for randomized data suggesting that for randomized sub-samples, the statistical expectations are confirmed. Despite shifts in the mean flow, the range of the expected error in standard deviation follows the correlation of two over the square root of the sample size minus one as suggested in [24]. The normalized probability density functions of the two velocity components at point A follow approximately a normal distribution. The mean was subtracted from \(V_x\) and \(V_z\) and each component was divided by its standard deviation in order to compare both in the same plot to the standard normal distribution (mean = 0, standard deviation = 1).

To illustrate the statistical error of the measurements as a function of crank angle, Figure 8 plots both the coefficient of variation \((COV = \text{spatial average of ensemble RMS divided by spatial average of ensemble average})\) and the statistical
error (COV divided by square root of sample size) at each crank angle. For most of the cycle the estimated average-velocity error is below 10%, with the exception of both top dead centers where an error of 18% can be expected. These errors are always less than 7% of the COV, which is a measure of CCV.

2 1D AND LARGE-EDDY SIMULATIONS

2.1 LES Methods and Models

LES at an engineering-level were conducted with the CONVERGE code [25] at GM R&D. Here, the computational domain includes the intake port and plenum, the cylinder, and exhaust port and plenum. 1D simulation results are used to define the initial and boundary conditions. Crank angle based total pressure and temperature are applied as inlet boundary conditions. Crank angle based static pressure is used as the outlet boundary condition. An orthogonal structured finite volume method with boundary cut-cell mesh and pressure-based time-implicit compressible schemes was employed, yielding second-order momentum and first-order space and time accuracies. Turbulent kinetic energy at the walls is replaced by sub-grid kinetic energy, approximated based on the sub-grid velocity as suggested in [26]. The Werner and Wengle wall model was used for the treatment of the sub-grid kinetic energy at the walls [27]. The mesh size in the near wall region is in the range of $y^+ = 10$ to 110.
Figure 7
Sample-size convergence of ensemble average velocity and of standard deviation. All data are normalized to the average of a 3 000-cycle run.
a) Color maps of coefficient of variance (COV) at 100 and 300 CAD ATDCE (vectors represent ensemble mean flow field). b) Distribution of sub-sample normalized average velocity magnitudes at point A as function of sample size (cycles). c) Distribution of sub-sample normalized standard deviation of velocity magnitudes at point A as function of sample size (cycles). d) Probability distribution of velocity magnitude follows normal distribution (3035 samples); Temporal evolution of 20 cycle sub-sample averages of both velocity components.
For the present work, two LES sub-grid scale turbulence models were used:

- one-equation eddy viscosity model [28, 29],
- Dynamic Structure One-Equation model (DSOE) [30, 31].

### 2.2 One-Equation Eddy Viscosity Model

The one-equation eddy viscosity model closes the sub-grid stress tensor with a turbulent viscosity. Model constants are required for the turbulent viscosity and the dissipation term. The model adds a transport equation for the sub-grid kinetic energy to the energy equation. The sub-grid kinetic energy is then used in modeling the turbulent viscosity and is given by:

\[
\frac{\partial k}{\partial t} + \mathbf{u}_i \frac{\partial k}{\partial x_i} = -\tau_{ij} \frac{\partial \mathbf{u}_i}{\partial x_j} - \epsilon + \frac{\partial}{\partial x_i} \left( \nu_i \frac{\partial}{\partial x_i} \right)
\]  \tag{1}

Here, the sub-grid kinetic energy is determined as:

\[
k = \frac{1}{2} \left( \mathbf{u}_i \mathbf{u}_i - \bar{\mathbf{u}}_i \bar{\mathbf{u}}_i \right)
\]  \tag{2}

The turbulent viscosity for the one-equation model is given as:

\[
\nu_i = C_k k^{1/2} \Delta
\]  \tag{3}

where \( \Delta \) the grid filter which is set to the local cell size. Note that the turbulent viscosity is a function of grid size so that in the limit of fine resolution the turbulent viscosity goes to zero. The turbulent viscosity can be ‘tuned’ by adjusting the constant \( C_k \) in the above expression and a value of 0.05 was used for the present work. The sub-grid dissipation is given as:

\[
\epsilon = C_v k^{3/2} / \Delta
\]  \tag{4}

The sub-grid dissipation can also be ‘tuned’ by adjusting the constant \( C_v \) in the above expression. Here, the value was set to 1.

### 2.3 Dynamic Structure One-Equation Model

The DSOE model is also a kind of one-equation LES model. The model for the sub-grid stress tensor does not use a turbulent viscosity concept but rather an attempt is made to model the sub-grid stress tensor more accurately (i.e., the sub-grid stress tensor cannot be accurately modeled with a turbulent viscosity). A dynamic approach is used to eliminate user-specified constants in the sub-grid stress tensor model. However, a dissipation model is still required that has a user-specified model constant.

To formulate a dynamic model, a second filtering operation, which is designated the test level filter, \( \Delta_t \), is required. The test filter, \( \Delta_t \), is typically twice the value of the grid filter, \( \Delta \). The residual stresses based on single (grid) and double filtering (test) operations are defined by:

\[
\tau_{ij} = \left( \mathbf{u}_i \mathbf{u}_j - \bar{\mathbf{u}}_i \bar{\mathbf{u}}_j \right)
\]  \tag{5}

and

\[
T_{ij} = \left( \mathbf{u}_i \mathbf{u}_j - \bar{\mathbf{u}}_i \bar{\mathbf{u}}_j \right)
\]  \tag{6}

The grid level stress tensor and the test level stress tensor are related by the Germano identity [32]:

\[
L_{ij} = T_{ij} - \tau_{ij} = \left( \bar{\mathbf{u}}_i \bar{\mathbf{u}}_j - \bar{\mathbf{u}}_i \bar{\mathbf{u}}_j \right)
\]  \tag{7}

Then the model for the sub-grid stress tensor becomes:

\[
\tau_{ij} = 2k \left( \frac{L_{ij}}{L_{ii}} \right)
\]  \tag{8}
For improved fidelity of wall flows simulation when running the LES model with CONVERGE, the Werner and Wengle wall model [33] was enabled.

Orthogonal cubic meshes were automatically created at run time to avoid any variation in cell shape or size. Adaptive Mesh Refinement (AMR) can also be used to increase the effective grid resolution but was not enabled in the present study. The LES sub-grid stress tensor modeled in CONVERGE yields 2nd order accuracy due to its central difference scheme.

A feature of the code called ‘embed-sphere’ was employed in all the LES calculations carried out at GM R&D, mainly to minimize the effect of mesh topology on the computed results and to enable better quantitative comparisons between various applications. An example is shown in Figure 2b where the finest mesh (0.5 mm) is placed in the valve seat and spark plug regions while the coarsest mesh (8 mm) is placed inside the intake and exhaust plenums.

3 RESULTS

The results of the experimental effort provide a set of benchmark data for motored flows to guide setup and eventually allow validation of multi-cycle LES calculations. Having demonstrated the test-to-test repeatability of the engine operation and data, this section demonstrates one application of the measured data for LES development. It is important to assess simulation performance for a range of operating conditions to avoid optimization or tuning of simulation settings to one condition but not gaining insight into the fidelity of reproducing different conditions. Two operating conditions will be highlighted here to illustrate how changes in intake manifold pressure affect the filling dynamics of the cylinder and how the simulations pick up this process. Conceptually, it is expected that a dynamic structure-based sub-grid treatment would perform better than a simple eddy viscosity model in handling the flow around the valves and the resulting strong velocity gradients in the intake ‘jet’ flow. Encouraging results that point to this then also lead to investigations of increasing the mesh resolution around the valves to better capture the dynamics of the valve opening and closing events; this resulted in higher accuracy for the numerical nature of valving events in CFD. In this case study comparisons are made between the 1 300 rpm/95 kPa and 1 300 rpm/40 kPa operating conditions. Prior to comparing the in-cylinder velocity distributions, it is of value to compare the mass-flow and pressures between the measured and simulated data.

3.1 Pressure Dynamics and Mass Flow

The TCC-III intake system is closed, meaning that the intake is not open to the atmospheric pressure. Instead, the critical orifice flow meters deliver a fixed mass flow rate, \( \dot{m} \), to the intake system and are adjusted to achieve the desire MAP. Thus, the intake plenum pressure fluctuates, since the plenum fills at fixed rate but empties according to the periodic flow rates of the engine’s intake stroke. The pressure was measured at the intake plenum inlet and exhaust plenum outlet and modeled quite accurately by the 1-D simulation as was shown in Figure 3a. Table A2 shows that the measured average mass flow is within 2.4% of that of the 1-D simulation. LES results obtained with GT Power model boundary conditions and experimental boundary conditions do not show significant differences. Thus, the 1-D model results were used as the crank-angle dependent plenum-boundary conditions for the LES calculations. In general, this enables simulations for operating conditions or engine designs for which no experimental data are available.

Since the plenum inlet/outlet pressures are the inputs set at run time, only the intake port, exhaust port and in-cylinder pressures are metrics of the LES effectiveness. As with the experimental test-to-test variation, Figure 3b quantifies the discrepancy of the phase-averaged pressures between the simulations (LES and 1-D) and measurements. Although the discrepancy is large compared to the test variability (Fig. 3a), the simulated intake port pressures are shown to be within 3% of the measured pressures during the intake stroke. The cylinder-pressure discrepancy is as much as 5% of the absolute values, but this is expected and acceptable due to the large dynamic range of the pressure transducer data.

The peak pressure and location of peak pressure are compared in Table A2 and shown to be within 7% and 1 CAD for the GT Power simulation and 2.5% and 1 CAD for LES. Both differences are dominated by the trapped mass, although small differences can be due to the heat transfer (Woschni model) and blowby, which was not modeled here.

Figure 9a shows the measured pressure drop across the intake port, valves, and the in-cylinder volume, \( \Delta P = P_{\text{Intake Port}} - P_{\text{Cylinder}} \). The simulated 1-D mass flow through the valve, \( \dot{m} \), and the valve lift are shown during the intake stroke (0-300 aTDCexh). The very good agreement between the 1-D simulation and measured \( \Delta P \) (refer to dashed-line, Fig. 9b), provides confidence that the 1-D simulated mass flow is reasonably accurate. Figure 9a shows high intake mass flow rates with relatively low lift before CAD 60, and negative flows during valve overlap (esp. 40 kPa, CA < 30) and between BDC and IVC. The importance of the mass-flow and lift during the early intake stroke will be discussed during the presentation of the velocity data.
3.2 Velocity Data

As an extension of the previous work at 800 rpm, 95 kPa [7], a goal of the present study was to develop the know-how of LES setup (sub-models/mesh refinement/valve timings etc.) pertaining to varying engine operating conditions of more practical relevance. The variation of intake manifold pressure is addressed as an example and explores predictive capabilities for 1 300 rpm, 95 kPa and 1 300 rpm, 40 kPa in comparison to comprehensive PIV data taken in four cross sections (Fig. 4). LES predictions with One-Equation Eddy Viscosity and One-Equation Dynamic Structure models are compared to PIV measurements for various crank angle positions. For illustration purposes, results are described at 100 CA, corresponding to the intake jet flow, and 300 CA, corresponding to the compression flow. Cutting planes in all three orthogonal directions (Y = 0, X = 0 and Z = −5 and −30 mm) are considered to get a three-dimensional view of the flow field, though only results of the horizontal plane at Z = −30 mm are discussed here. After a first look into the 1 300 rpm, 95 kPa case and the lower MAP operating point (1 300 rpm, 40 kPa) is explored to identify differences in the flow that are expected from the observed pressure dynamics (Fig. 9).

Ensemble average velocity vector fields and ensemble RMS fields were visually compared side-by-side. Only the in-plane velocity components were considered for LES, because the PIV measurements provide only two velocity components. At every point in each plane, the ensemble average and ensemble RMS were calculated, over all cycles n, as follows in m/s:

\[
V_{avg} = \frac{1}{n} \sum_{i=1}^{n} V_i 
\]

(9)

\[
V_{rms} = \sqrt{\frac{1}{n-1} \sum_{i=1}^{n} (V_i - V_{avg})^2}
\]

(10)

Figure 9

Comparison between the measured and simulated pressure drop across the port and valve (\(\Delta P = P_{\text{IntakePort}} - P_{\text{cyl}}\)). 1-D mass flow through the valve, and valve lift during the intake stroke; Note that in a) the bottom part of the y-axis is a continuation of \(\Delta P\) for valve opening and closing.

a) Measured \(\Delta P\), 1-D simulation of mass flow, and valve lift. b) Discrepancy between the measured and simulated \(\Delta P\).
Figure 10

PIV and LES ensemble averaged mean flow velocities (velocity magnitude color-coded, in m/s) at various cutting planes and crank angle positions for 1 300 rpm, 95 kPa. The circles indicate the PIV FOV. Number of cycles used for ensemble average calculation: PIV $Y = 0$ mm $\# = 240$, PIV $Z = -30$ mm $\# = 1,157$, one-equation $\# = 20$, dynamic structure $\# = 20$. 
Figure 11

PIV and LES ensemble averaged RMS velocities (RMS velocity magnitude color-coded, in m/s) at various cutting planes and crank angle positions for 1 300 rpm, 95 kPa. Number of cycles used for RMS calculation is the same as in Figure 10.
Previous work [7] showed good agreement between LES one-equation model and PIV for 800 rpm, 95 kPa. Thus, a first comparison will be shown for results from using the one-equation model with PIV for the current operating point.

The mean flow predicted by the one-equation model is shown in Figure 10. Intake jet penetration depth (100 CA, $Y = 0$ plane) seems to be a little under-predicted. In the azimuthal (100 CA, $Z = -30$) cutting plane, the flow structures are captured well qualitatively, but the overall magnitude is smaller than in the experiments. During the compression stroke at 300 CA, for all cutting planes the mean velocity fields match reasonably well qualitatively, but the predicted magnitude is consistently lower than the experimental data.

Next the performance of the DSOE model was examined, which is expected to better resolve the sub-grid kinetic

---

**Figure 12**

PIV and LES ensemble averaged mean flow velocities (velocity magnitude color-coded, in m/s) at 100 deg aTDCE and various cutting planes for 1 300 rpm, 40 kPa. Number of cycles used for ensemble average calculation: PIV $Y = 0$ mm # = 240, PIV $X = 0$ mm # = 235, PIV $Z = -30$ mm # = 1 157, one-equation # = 20, dynamic structure # = 20.
energy because of a dynamic interaction with the resolved kinetic energy, rather than using the same model constant to compute the sub-grid kinetic energy as in the one-equation model. Figure 10 shows that there is not much improvement with the DSOE model over the one-equation model in terms of intake jet shape or penetration depth ($Y = 0$ plane), but the agreement in the azimuthal ($Z = -30$ plane) becomes much better at 100 CA for the former. Interestingly, the DSOE model under-predicts the compression flow ($Y = 0$ plane at 300 CA) and the overall magnitude is even lower than for the one-equation model. There is not much difference in the two models for azimuthal ($Z = -30$ plane) mean flow at 300 CA beyond a small difference in the swirl magnitude.

Figure 11 shows that RMS is consistently under-predicted especially during the intake stroke. However, some regions of large RMS (within and at the bottom of the intake jet at 100 CA, $Y = 0$ plane; the bottom of the domain at 300 CA, $Y = 0$ plane and $Z = -30$ plane) are correctly captured by both the models. At 100 CA, $Z = -30$ plane,
neither model captures the correct qualitative trends. Our previous study [7] shows that by adding the SGS turbulent fluctuation $u''$ to LES-calculated RMS values, the agreement between PIV measurement and LES prediction can be further improved.

From the above discussion, it seems that both one-equation and DSOE models do a good job in capturing some of the features of the flow field, but no overarching conclusions as to which model performs better can be made for this operating point.

The overall goal of the GM-UM collaboration is to develop LES tools to predict cyclic variability of engine flows in the presence of combustion. For optical engine experiments with combustion, the intake manifold pressure has to be maintained at a much lower level than the ambient because the optical cylinder setup is rated only to 6 MPa peak pressure. In this study therefore the motored flow at a MAP of 40 kPa was examined, as a pre-cursor to capturing flows with combustion.

**Figure 14**
PIV and LES ensemble RMS flow velocities (RMS velocity magnitude color-coded, in m/s) at 100 deg aTDCE and various cutting planes for 1 300 rpm, 40 kPa. Number of cycles used for RMS calculation is the same as in Figure 12.
As in the previous section, both one-equation and DSOE models were studied with a compilation of results shown in Figures 12-15. The velocity fields shown are taken for the same CA positions as before (intake jet: 100 CA and compression flow: 300 CA). Clearly, the use of the DSOE model results in a prediction that is far closer to the experimentally determined flow fields at 100 CA than what is obtained with the one-equation model (Fig. 12). The shape of the intake jet along with its penetration depth and the surrounding vortical flow structures are captured more accurately in the $Y = 0$ plane.

The $X = 0$ cross section of the intake jet follows that trend as well. For the $Z = -30$ plane the magnitude and overall structural features of the flow are again better captured with the DSOE model. Figure 13 shows that at 300 CA, however, both the models predict similar trends for all cutting planes.

In general, RMS magnitudes are consistently under-predicted by both the models (Fig. 14, 15), with the DSOE model predicting more accurate RMS distributions than...
one-equation model at 100 CA for the $X$ and $Y = 0$ planes. However, in the horizontal plane and at 300 CA both the models predict inaccurate distributions.

To summarize, for the 1 300 rpm, 40 kPa operating point, both models give similar results for compression flow (300 CA), but the DSOE model performs better than one-equation model in capturing intake flow (100 CA), both in terms of mean and RMS velocity distributions. This is an important finding, as it is contrary to what was observed for 95 kPa, where both models performed similarly for both intake and compression flows.

The reason for this unexpected behavior could be found in the differences of the bulk intake flow that was discussed in the context of Figure 9. What is the difference between the intake flow at 95 kPa versus 40 kPa? As seen in Figure 9a, there is a larger backflow from cylinder to intake port for 40 kPa because of a reverse pressure gradient. This makes the overall intake flow structure at 40 kPa much more complicated (in-flow and out-flow through the intake valve) than at 95 kPa. This implies that for such complex flows, the one-equation model fails to provide accurate results, and one needs to resort to the DSOE model.
In order to further substantiate the above point, the early intake flow was examined in greater detail. Figures 16 and 17 show the mean and RMS, respectively at 60 CA. Clearly, the DSOE model predicts the downward curved intake jet shape ($Y = 0$ plane) consistent with PIV, as opposed to a much weaker intake jet predicted by the one-equation model. Flow structures at $X = 0$ plane and $Z = -30$ plane are similar between the two models, though. High RMS regions are adequately captured by both the models for all the cutting planes, although here too, the DSOE model performs marginally better.

**Mesh Refinement**

The observation that the early phase of the intake flow is so critical to the evolution of the flow indicated that the DSOE model could do better in capturing the high-gradient
flow regions. Another critical aspect, though, is the way valve opening and closing events are handled in CFD. In the physical engine setup and the experiments the valve lift curves are continuous and smooth functions in time. In a CFD simulation, however, the valve opening event is an abrupt process that will switch the valve position from fully closed to open at a pre-set valve lift, typically a fraction of a millimeter, from one time step to the next. Thus, the initial flow conditions will be substantially affected by this minimum valve lift setting, especially for conditions where there is a substantial initial pressure difference across the valve.

A potential means to better capture the complex flow at 40 kPa then is to refine the mesh close to the valve seats, as finer meshes may allow resolving the detailed flow structures better. In an attempt to improve the earlier one-equation model results for intake flow at 40 kPa, the smallest mesh size close to the valve seat was reduced to 0.25 mm as opposed to the regular 0.5 mm. This comes at a substantial computational expense and the simulations with finer mesh took three times longer to finish.

The mean flow results are presented in Figure 18 for 20 cycles and clearly show that mesh refinement leads to much better results; with more accurate intake jet penetration depth in the \( Y = 0 \) plane, and better quantitative agreements at the bottom of the spark plug in the \( X = 0 \) and the \( Z = -30 \) planes. For comparison, the results from the DSOE model but with coarser grid are also shown again here. This illustrates how mesh refinement and the increased ability of a DSOE model better capture gradients. Figure 19 makes the same point for the RMS velocity distributions. Thus, it is expected that future computations with a finer mesh around the valves and the DSOE model will further increase the accuracy of the simulations.

**Figure 18**

Mesh refinement study: PIV and LES ensemble-averaged mean flow velocities (velocity magnitude color-coded, in m/s) at 100 deg aTDCE and various cutting planes for 1 300 rpm, 40 kPa. Number of cycles used for ensemble average calculation: PIV \( Y = 0 \) mm \# = 240, PIV \( X = 0 \) mm \# = 235, PIV \( Z = -30 \) mm \# = 1 157, one-equation (0.5 mm) \# = 20, one-equation (0.25 mm) \# = 20, dynamic structure \# = 20.
Further experimental results are obtained at 800 RPM and 95 kPa in the TCC-III configuration (not presented here, but are available for download from http://deepblue.lib.umich.edu/handle/2027.42/108382) with the purpose to inspire future investigations that explore the impact of SGS models and meshing details by using parametric variations.

CONCLUSIONS

A comprehensive study was conducted to characterize the flow field in an optical engine, the TCC-III engine, to advance the predictive capability of LES. Experiments, 1D simulations, and LES were performed for motored engine conditions as a baseline for a study of the origins of cyclic variability in engine performance. High-speed PIV was used to capture the flow in hundreds of consecutive cycles with 5 CA degree steps. Successive measurements in all three orthogonal directions were performed over the course of a six-month period. Repeatability from test to test and within individual runs was shown to be very high, exceeding the experimental uncertainty of the measurements. The experiments were carried out for 1300 rpm and 40 and 95 kPa intake manifold pressures as well as for 800 rpm and 95 kPa MAP to enable a thorough assessment of the performance of two subgrid scale models used in LES that were performed with CONVERGE. A one-equation eddy viscosity model and a DSOE model were used with CONVERGE to perform computations for 20 consecutive engine cycles for each operating condition.

The experimental data was used to guide the selection of SGS models and assess their strength and weaknesses. In that context, the investigation of multiple operating conditions was critical for the analysis.

The simulation results from both SGS models show a trend that under predicts the magnitude of the mean flow.
and of the RMS velocity fields. However, it is noted that the spatial structure of the RMS velocity field in some observation planes is better captured in LES using the DSOE model. The differences were highest during the intake stroke and less pronounced during the compression stroke. This is plausible because the performance of the DSOE model is expected to be better in regions with high velocity gradients, such as during the intake when fast flowing air enters the cylinder. Contrasting the flows observed in low MAP operation (40 kPa) with those seen at 95 kPa reveals substantial differences in the dynamics of the bulk flow.

An investigation of early stage flow structures after intake valve opening pointed to the limited resolution of the computational mesh and the stepwise opening of the valve, a necessary procedure in CFD. Refined mesh of computational mesh from 0.5 mm to 0.25 mm around the valve improved the simulation accuracy even for the one-equation model quite substantially. Future computations will include refined meshing around the valve in combination with the DSOE model and also a parametric study of the flow at different engine speeds.

The data, along with the geometry of the engine and a complete GT Power model are available for benchmarking and model development from:
http://deepblue.lib.umich.edu/handle/2027.42/108382

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18 Abraham P., Liu K., Haworth D., Reuss D., Sick V. (2014) Evaluating Large-Eddy Simulation (LES) and High-Speed Particle Image Velocimetry (PIV) with Phase-Invariant Proper Orthogonal Decomposition (POD), Oil & Gas Science and Technology 69, 1, 41-59.


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APPENDIX A

TCC-III Engine and Data Specifications

<table>
<thead>
<tr>
<th>TABLE A1</th>
<th>Engine geometry</th>
</tr>
</thead>
<tbody>
<tr>
<td>Bore (cm)</td>
<td>9.20</td>
</tr>
<tr>
<td>Stroke (cm)</td>
<td>8.60</td>
</tr>
<tr>
<td>Clearance @ TDC</td>
<td>0.95</td>
</tr>
<tr>
<td>Comb chamber (cc)</td>
<td>63.15</td>
</tr>
<tr>
<td>Top land crevice (cc)</td>
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</tr>
<tr>
<td>Spkplug crevice (cc)</td>
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<tr>
<td>TDC vol. (cc)</td>
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<tr>
<td>Swept vol. (cc)</td>
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</tr>
<tr>
<td>Geometric CR</td>
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</tr>
<tr>
<td>Effective (IVC) CR</td>
<td>8.00</td>
</tr>
<tr>
<td>Steady-flow swirl ratio</td>
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</table>

<table>
<thead>
<tr>
<th>TABLE A2</th>
<th>Data precision</th>
</tr>
</thead>
<tbody>
<tr>
<td>Set point parameter name</td>
<td>Set point</td>
</tr>
<tr>
<td>Engine speed (RPM)</td>
<td>800/1300/1300</td>
</tr>
<tr>
<td>Pressure at intake plenum inlet (kPaa)</td>
<td>95/95/40</td>
</tr>
<tr>
<td>Pressure at exhaust plenum outlet (kPaa)</td>
<td>101.5/101.5/101.5</td>
</tr>
<tr>
<td>Air temperature intake port (°C)</td>
<td>45/45/45</td>
</tr>
<tr>
<td>Engine oil temperature at inlet (°C)</td>
<td>45/45/45</td>
</tr>
<tr>
<td>Engine coolant temperature at outlet (°C)</td>
<td>45/45/45</td>
</tr>
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</table>

(continued)
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<tr>
<th>Measured parameter name</th>
<th>Average of all tests</th>
<th>Max deviation from all test avg to single test avg</th>
<th>Average in test stdev</th>
<th>GT power model</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pressure at intake port (kPaa)</td>
<td>94.6/94.4/39.5</td>
<td>0.4/0.1/0.3</td>
<td>0.05/0.1/0.05</td>
<td>95.1/94.8/39.8</td>
</tr>
<tr>
<td>Pressure at exhaust port (kPaa)</td>
<td>101.7/101.8/101.1</td>
<td>0.1/0.1/0.1</td>
<td>0.03/0.01/0.03</td>
<td>101.2/101.1/100.9</td>
</tr>
<tr>
<td>Indicated mean effective pressure (kPa)</td>
<td>–41.9/–37.8/–19.6</td>
<td>0.9/0.1/1.0</td>
<td>0.7/0.5/0.3</td>
<td></td>
</tr>
<tr>
<td>Cylinder peak pressure (kPaa)</td>
<td>1810/1958/798</td>
<td>16/0.6/7</td>
<td>3.4/3.5/1.8</td>
<td>1937/2013/836</td>
</tr>
<tr>
<td>Cylinder peak pressure location (CA aTDCE)</td>
<td>358.6/358.8/358.5</td>
<td>0.01/0.03/0.07</td>
<td>0.2/0.25/0.3</td>
<td>359.5/359.5/359.5</td>
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<td>Temperature at outside cylinder wall (°C)</td>
<td>36.9/40.8/41.0</td>
<td>0.5/2.6/0.3</td>
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<tr>
<td>Mass flow (g/s)</td>
<td>3.5/5.88/2.07</td>
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<td>0/0/0</td>
<td>3.51/5.80/2.12</td>
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</table>

**LES model results**

<table>
<thead>
<tr>
<th>Computed parameter</th>
<th>800 95 OE</th>
<th>1300 95 OE</th>
<th>1300 40 OE</th>
<th>1300 40 OE refined mesh</th>
<th>130040 DS</th>
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</thead>
<tbody>
<tr>
<td>Pressure at intake port (kPaa)</td>
<td>94.8</td>
<td>94.5</td>
<td>39.7</td>
<td>39.8</td>
<td>39.6</td>
</tr>
<tr>
<td>Pressure at exhaust port (kPaa)</td>
<td>101.4</td>
<td>101.3</td>
<td>100.8</td>
<td>100.8</td>
<td>100.8</td>
</tr>
<tr>
<td>Cylinder peak pressure (kPaa)</td>
<td>1863</td>
<td>1941</td>
<td>807</td>
<td>810</td>
<td>813</td>
</tr>
<tr>
<td>Cylinder peak pressure location (CA aTDCE)</td>
<td>359.5</td>
<td>359.5</td>
<td>359.5</td>
<td>359.3</td>
<td>359.5</td>
</tr>
<tr>
<td>Mass flow (g/s)</td>
<td>3.50</td>
<td>5.85</td>
<td>2.13</td>
<td>2.12</td>
<td>2.13</td>
</tr>
</tbody>
</table>

Note: Motored parameter values in the table are organized according to engine speed and MAP operating condition, i.e., parameter value at (800 rpm, 95 kPa)/ (1 300 rpm, 95 kPa)/(1 300 rpm, 40 kPa).
<table>
<thead>
<tr>
<th>PIV plane</th>
<th>Y = 0 mm</th>
<th>X = 0 mm</th>
<th>z = −5 mm</th>
<th>z = −30 mm</th>
</tr>
</thead>
</table>
| 800 RPM 95 kPa | S_2014_05_20_02  
0-705CAD  
40-705CAD  
235cyc | No data | S_2014_05_13_03  
0-705CAD  
0-705CAD  
286cyc | S_2 014_04_17_02  
55-300CAD  
60-300CAD  
1157cyc |
| 1300 RPM 95 kPa | S_2014_05_20_01  
0-705CAD  
40-705CAD  
235cyc | No data | S_2014_05_13_02  
0-705CAD  
0-705CAD  
407cyc | S_2014_04_17_0  
155-300CAD  
60-300CAD  
1157cyc |
| 1300 RPM 40 kPa | S_2013_10_24_01  
0-690CAD  
0-690CAD  
240cyc | S_2014_01_30_01  
0-705CAD  
0-705CAD  
235cyc | S_2014_05_05_01  
0-705CAD  
0-705CAD  
407cyc | S_2014_03_26_01  
55-300CAD  
60-300CAD  
1134cyc |
|             | S_2013_11_07_02  
0-690CAD  
0-690CAD  
240cyc | S_2014_02_04_01  
0-705CAD  
0-705CAD  
235cyc | S_2014_05_05_02  
0-705CAD  
0-705CAD  
407cyc | S_2014_04_03_01  
55-300CAD  
60-300CAD  
1134cyc |
|             | S_2013_11_11_01  
0-712.5CAD  
0-712.5CAD  
116cyc | S_2014_02_05_02  
0-705CAD  
0-705CAD  
235cyc | S_2014_05_07_01  
0-705CAD  
0-705CAD  
407cyc | S_2014_04_16_02  
55-300CAD  
60-300CAD  
1157cyc |
|             | S_2014_05_20_03  
40-300CAD  
50, 60, 70, 80, 90, 100, 180, 220, 260, 300CAD  
3035cyc | | | S_2014_04_16_03_A  
55-300CAD  
60-300CAD  
400cyc |